

Nonlocal control of the Kondo effect in a double quantum dot–quantum wire coupled system

S. Sasaki,^{1,*} S. Kang,^{2,3} K. Kitagawa,^{4,3} M. Yamaguchi,^{1,3} S. Miyashita,⁵ T. Maruyama,⁵ H. Tamura,^{1,3}
T. Akazaki,^{1,3} Y. Hirayama,^{1,6} and H. Takayanagi^{1,4,3}

¹*NTT Basic Research Laboratories, NTT Corporation, Atsugi, Kanagawa 243-0198, Japan*

²*Institute of Physics, University of Tsukuba, Tsukuba, Ibaraki 305-8571, Japan*

³*CREST-JST, Honmachi, Kawaguchi, Saitama 331-0012, Japan*

⁴*Faculty of Science, Tokyo University of Science, Kagurazaka, Tokyo 162-8601, Japan*

⁵*NTT Advanced Technology Corporation, Atsugi, Kanagawa 243-0198, Japan*

⁶*SORST-JST, Honmachi, Kawaguchi, Saitama 331-0012, Japan*

(Received 22 March 2006; published 20 April 2006)

We have performed low-temperature transport measurements on a double quantum dot–quantum wire coupled device and demonstrated nonlocal control of the Kondo effect in one dot by manipulating the electronic spin states of the other. We discuss the modulation of the local density of states in the wire region due to the Fano-Kondo antiresonance, and the Ruderman-Kittel-Kasuya-Yoshida exchange interaction as the mechanisms responsible for the observed features.

DOI: [10.1103/PhysRevB.73.161303](https://doi.org/10.1103/PhysRevB.73.161303)

PACS number(s): 73.23.Hk, 73.63.Kv, 75.20.Hr

Semiconductor quantum dots (QDs) have attracted growing interest over the past years because of their controllable electronic states, and are often referred to as artificial atoms. One aspect of the QD as a single magnetic impurity culminated in the realization of an artificial Kondo system.^{1–6} When the QD has an odd number, N , of electrons, it normally has spin $S=1/2$. Then, the Kondo effect occurs when the temperature is below the Kondo temperature, T_K , and manifests itself as an increase in the conductance at Coulomb blockade regions and as a zero-bias peak in the differential conductance versus source-drain bias characteristics. These features are observed for the conventional geometry where the source and drain electrodes are attached to both ends of the QD, i.e., the only current path is through the “magnetic impurity” (embedded geometry). On the other hand, one can realize side-coupled geometry where a QD is attached to the side of another quantum wire (QW) without direct coupling to a source or drain electrode. In such a case, the Fano resonance appears in the conductance of the QW due to the interference between a localized state in the QD and the continuum in the QW (Refs. 7 and 8). When the QD has odd N , the Kondo “cloud” suppresses the conductance of the QW in the region between two Fano resonances, which is referred to as the Fano-Kondo antiresonance.^{9–12} An interesting spin correlation can be investigated when two QDs, each carrying spin $1/2$, are coupled to the QW. In this case, indirect exchange interaction is expected between the two localized spins mediated by conduction electrons in the QW, which is known as the Ruderman-Kittel-Kasuya-Yoshida (RKKY) interaction.^{13–18} The RKKY interaction may also be of great importance in the context of quantum information processing as a means to correlate remote spin quantum bits connected to a common conduction channel.

In this Rapid Communication, we present transport measurement on two QDs coupled to the opposing sides of a quasi-one-dimensional QW. This system offers smaller interdot distance in nominal device geometry than the previous experiment by Craig *et al.*,¹⁵ and is advantageous for studying the RKKY interaction, which decays on the length scale

of the Fermi wavelength. Moreover, the “effective” interdot distance can be tuned to zero in an ideal case of well-defined QW and zero relative distance along the wire direction.^{13,14} We demonstrate nonlocal control of the Kondo effect in one dot by tuning the spin states of the other, which suggests the RKKY interaction as an underlying mechanism.

The inset of Fig. 1(a) shows the double QD-QW coupled device fabricated by depositing Schottky gates on a GaAs/Al_xGa_{1-x}As heterostructure having two-dimensional electron gas 90 nm below the surface. Two QDs and a QW between them are formed (shown by dotted lines) by biasing the eight gates. To avoid formation of another QD in the central QW region, V_{sR} and V_{sL} are biased more negatively than V_{mR} and V_{mL} . Transport measurement is performed with a standard lock-in technique at temperature $T \sim 100$ mK in a dilution refrigerator. The source-drain bias $V_{sd} = V_{dc} + V_{ac}$, where V_{dc} is a dc offset, and ac excitation (about 13 Hz) $V_{ac} = 3 \mu\text{V}$.

Figure 1(b) shows a color-scale plot of the linear conductance through the right QD as a function of V_{pL} and V_{pR} . Attention is paid to completely pinch off the exit tunnel barrier of the left QD by applying sufficiently negative voltage, V_{dL} . Therefore, the only current path is through the right QD. This is done to avoid a spurious mirror effect, which occurs in the presence of multiple current paths.¹⁹ The gate voltage that controls the coupling of the left QD with the rest of the system, $V_{sL} = -1.04$ V, is adjusted in the weak-coupling regime. The two dotted lines designate Coulomb oscillation peaks for the right QD, and the region between them is a Kondo valley where the number of electrons in the right QD, N_R , is odd. The conductance in the Kondo valley is enhanced to almost the same value as that at the Coulomb oscillation peaks, indicating that the measurement temperature is well below T_K . Conductance maxima appear, forming almost vertical ridges as V_{pL} is swept. When we monitor the current through the QW only, I_W , instead of that through the right QD, I_R , as a function of V_{pL} , the Fano resonance associated with the left QD is observed as shown in Fig. 1(a). Here, the Fano’s asymmetric parameter, q , seems to be large, produc-

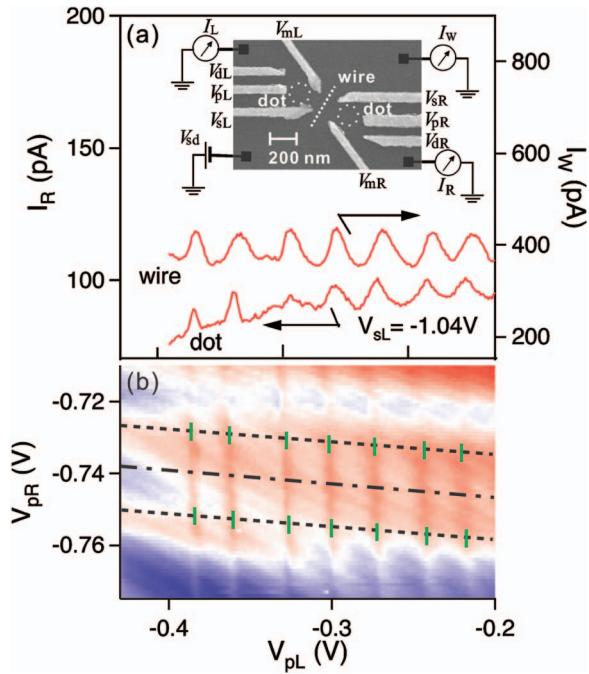


FIG. 1. (Color) (a) The current through the right QD, I_R , and that through the QW, I_W , in response to excitation source-drain voltage $V_{ac}=3 \mu\text{V}$ as a function of V_{pL} . V_{pR} is changed at the same time to trace the dash-dotted line shown in Fig. 1(b). I_R and I_W are measured separately, but with the same gate voltage conditions, which determine the wire conductance to be around $4e^2/h$. The inset shows a scanning electron micrograph of the double QD-QW coupled device together with a schematic of the measurement setup. Dotted lines indicate the positions of the QDs and QW. (b) Color-scale plot of the conductance through the right QD as a function of V_{pL} and V_{pR} in the weak-coupling regime with $V_{sL}=-1.04 \text{ V}$. Blue corresponds to $10 \mu\text{S}$, white to $26 \mu\text{S}$, and red to $42 \mu\text{S}$. Dotted lines denote Coulomb oscillation peaks for the right QD, and green vertical bars denote resonances of the left QD states with the chemical potential of the lead. Faint diagonal lines originate from the measurement setup.

ing almost symmetric peaks at the resonances for the present conditions.²⁰ I_R and I_W exhibit almost in-phase oscillations with the period of 20–30 mV in V_{pL} . This period is also consistent with the separation between Coulomb oscillation peaks observed in the direct current through the left QD, I_L , which flows with more positive V_{dL} . Therefore, we assign the ridges shown in Fig. 1(b) to resonances of the left QD states with the chemical potential of the lead. The Fano resonance modulates the conductance of the QW, and I_R , which flows via the QW, is also modulated. The ridges are not observed when V_{sL} is smaller than about -1.2 V because the left QD is decoupled from the rest of the system. Although the Kondo valley conductance is somewhat reduced from the original value at $V_{pL}=0$, the Kondo effect still remains between the ridges. Please note that the two leftmost ridges are separated from the others by a larger gap, suggesting a spin pair where up-spin and down-spin electrons consecutively occupy the same orbital state. However, there is no noticeable difference in the conductance between even and odd N_L regions, where N_L is the number of electrons in the left QD.

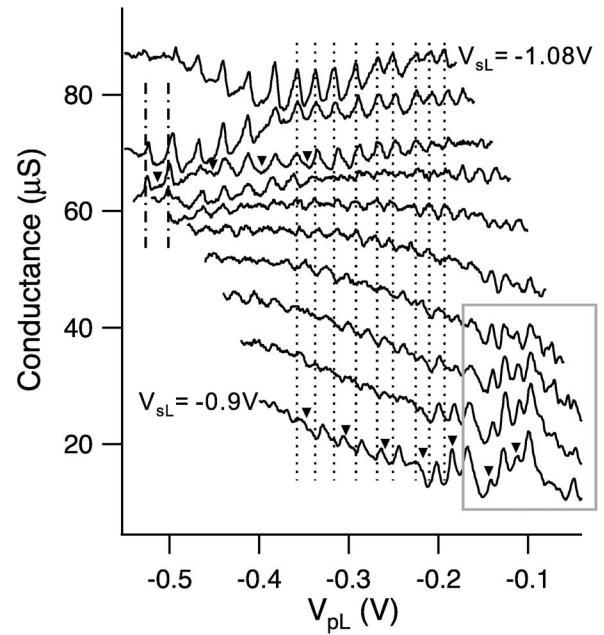


FIG. 2. Conductance profiles taken along the mid-Kondo valley of the right QD with the coupling gate voltage V_{sL} changed between -0.9 and -1.08 V in 20 mV steps. Curves are offset horizontally to compensate for the electrostatic energy shift due to V_{sL} and also vertically for clarity. The scales on the axes apply to the lowermost $V_{sL}=-0.9 \text{ V}$ trace. The dotted lines are guides for the eye, and the two dash-dotted lines on the left denote spin-pair peaks. Odd N_L regions are marked with solid triangles. The boxed region is plotted in more detail in Fig. 4.

Next, we increase the coupling of the left QD by increasing V_{sL} . Figure 2 shows evolution of the conductance profiles taken along the mid-Kondo valley of the right QD [similar to Fig. 1(a)] when V_{sL} is changed from -1.08 V (weak-coupling regime) to -0.9 V (strong-coupling regime). Asymmetric peaks reflecting the Fano resonances are observed in the weak-coupling regime. We estimated the shift in V_{pL} of the resonance peaks in the weak-coupling regime in response to a slight shift in V_{sL} that does not affect the Fano line shape, and found that the electrostatic energy shift caused by V_{sL} is almost identical to that caused by V_{pL} in our device. Based on this finding, we shift the traces in Fig. 2 horizontally so that the resonances arising from the same N_L states are aligned vertically. Using the dotted lines as a guide, one can see that the resonances appearing as conductance maxima in the weak-coupling regime change into minima in the strong-coupling regime, just as in the previous report by Sato *et al.*¹² We assign odd N_L regions and mark them with solid triangles, using the spin pair behavior in Fig. 1(b) as a reference (the two dashed lines in Fig. 2).

Figure 3(a) shows a plot similar to Fig. 1(b), but with $V_{sL}=-0.9 \text{ V}$, in the strong-coupling regime. Here, in contrast to Fig. 1(b), the Fano resonance involving the left QD with $q \approx 0$ produces conductance minima [green vertical bars in Fig. 3(a)]. The intensity of the several conductance maxima corresponding to fixed N_L seems to alternate between $V_{pL} = -130$ and -70 mV . As shown in Fig. 3(b), the zero-bias Kondo peak still remains in the dI_R/dV_{dc} versus V_{dc} charac-

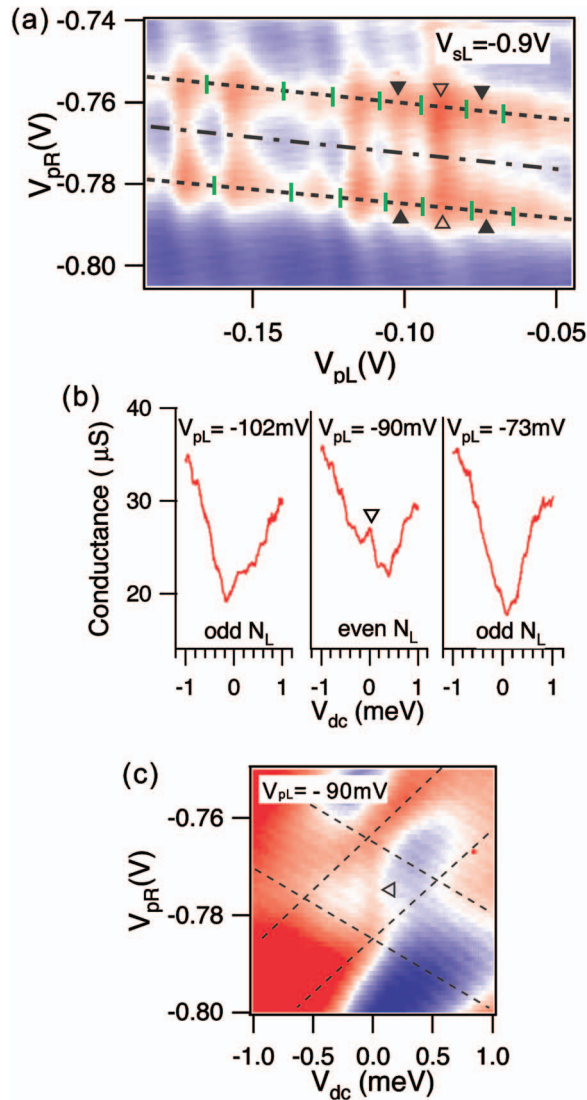


FIG. 3. (Color) (a) Color-scale plot of the conductance through the right QD as a function of V_{pL} and V_{pR} with $V_{sL} = -0.9$ V in the strong-coupling regime. Blue corresponds to $0 \mu\text{S}$, white to $18 \mu\text{S}$, and red to $36 \mu\text{S}$. Dotted lines and green vertical bars have the same meanings as in Fig. 1(b). The open and solid triangles denote even and odd N_L regions, respectively. (b) dI_R/dV_{dc} vs V_{dc} for odd N_L [$V_{pL} = -102$ and -73 mV, the solid triangles in (a)] and for even N_L [$V_{pL} = -90$ mV, the open triangle in (a)]. V_{pR} is fixed at -775 mV in the middle of the odd N_R Coulomb blockade valley. The triangle for $V_{pL} = -90$ mV case denotes a Kondo resonance peak. (c) Conductance of the right QD as a function of V_{dc} and V_{pR} for even N_L . Blue corresponds to $10 \mu\text{S}$, white to $25 \mu\text{S}$, and red to $40 \mu\text{S}$. The triangle denotes a Kondo ridge observed in the odd N_R Coulomb blockade region. The edges of the Coulomb diamonds are shown by dashed lines.

teristics when N_L is even ($V_{pL} = -90$ mV). The Kondo effect for this condition is also confirmed in the Coulomb diamond characteristics shown in Fig. 3(c). On the other hand, there is no Kondo peak when N_L is odd [$V_{pL} = -102$ and -73 mV, solid triangles in Fig. 3(a)]. Thus, a nonlocal control of the Kondo effect in the right QD is realized by changing the number of electrons in the left QD.

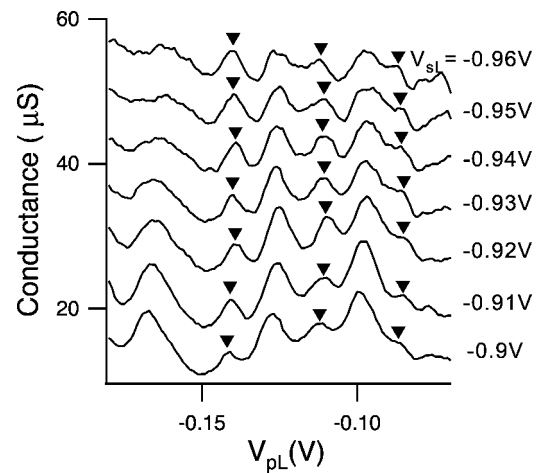


FIG. 4. Conductance profiles taken along the mid-Kondo valley of the right QD when the coupling of the left QD is changed with V_{sL} . Each curve is offset by $5 \mu\text{S}$ vertically and by -10 mV horizontally with respect to the previous curve with 10 mV more V_{sL} . The scales on the axes apply to the lowermost trace. The solid triangles denote odd N_L conductance peaks. There is an overall shift of about 10 mV in V_{pL} compared to Fig. 3 due to the different measurement time.

Figure 4 shows the V_{sL} dependence of the mid-Kondo-valley conductance profiles taken from the boxed region in Fig. 2 with smaller 10 mV steps in V_{sL} . Dips correspond to the Fano resonances associated with the left QD, and peaks correspond to the fixed N_L regions. The intensity of the odd N_L peaks (marked with triangles) is not very different from that of the even N_L peaks when the coupling is relatively weak ($V_{sL} \approx -0.96$ V). As V_{sL} is made more positive (≈ -0.92 V), dips due to the Fano resonance develop and make the peaks clearer, while the overall conductance slightly decreases. When V_{sL} is further increased and the left QD is strongly coupled, odd N_L peaks are suppressed more than even N_L ones, leading to the suppression of the Kondo effect as seen in Fig. 3(b).

Let us now consider the mechanisms responsible for the observed even-odd effect. As already mentioned, the dot current, I_R , is measured via the QW region, i.e., the QW acts as one lead for the right QD. When the Fano resonance involving the left QD modulates the conductance of the QW, it also indirectly modulates I_R . As coupling of the left QD with the QW increases, the Fano-Kondo antiresonance effect suppresses the conductance at the odd N_L regions,¹² consistent with the behavior observed in Fig. 4. It must be noted that the observed even-odd effect is different from a simple series resistance modulation. If that were the case, the zero-bias Kondo peak would still appear, but with a reduced height, in the case of odd N_L too. In reality, however, the Kondo temperature itself is changed between even and odd N_L . It is expected that the Fano resonance (or Fano-Kondo antiresonance) modulates the local density of states of the QW, as the enhanced zero-bias conductance shown in Fig. 3 in Ref. 21 suggests. Because $T_K \propto \exp(-1/\rho J)$, where ρ is the density of states and J is the coupling constant,²² T_K decreases when ρ decreases. This may be the case for the odd N_L peaks in Fig. 4 because ρ is reduced by the Fano-Kondo effect compared

to even N_L ones. Usually, the Kondo effect is experimentally controlled with a variety of parameters such as temperature, the tunnel coupling with the leads, the position of the localized level relative to the Fermi energy, source-drain bias, and magnetic field. The control of T_K via the Fano(-Kondo)-modulated density of states represents a totally unique way to experimentally switch on/off the Kondo effect. This Fano-Kondo antiresonance mechanism is in accordance with the recent theory by Martins *et al.*,²³ that gives an alternative explanation of the results by Craig *et al.*¹⁵

Another mechanism that could be responsible for the observed even-odd effect is the RKKY interaction. When both N_R and N_L are odd, the RKKY interaction couples two spins in the two QDs via interceding conduction electrons either ferromagnetically or antiferromagnetically, depending on the distance between them. The Kondo effect for the single QD with spin 1/2 is suppressed both in the former case (the total spin, $S_{\text{tot}}=1$), corresponding to the underscreened Kondo effect, and in the latter case ($S_{\text{tot}}=0$). The effective interdot distance in our device structure is considered to be close to zero since they couple to the opposing sides of the quasi-

one-dimensional QW (having two to three conduction channels for the gate voltage conditions explored) with a small relative distance along the wire direction.^{13,14} Hence, the ferromagnetic coupling may be realized, and relatively large interaction strength is expected. However, this RKKY mechanism is inseparable from the Fano-Kondo-modulated density-of-states scenario described above, and presumably both mechanisms contribute to the observed suppression of the Kondo effect when the remote QD also carries spin 1/2.

In summary, we have presented transport measurement on a double QD-QW coupled system and demonstrated nonlocal control of the Kondo effect in one QD by manipulating the spin states of the other. As for the responsible mechanisms, we consider modulation of the Kondo temperature by the modulated local density of states of the QW due to the Fano-Kondo antiresonance, and the RKKY interaction between the two QDs.

This work is financially supported by a Grant-in-Aid for Scientific Research from the Japan Society for the Promotion of Science.

*Electronic address: satoshi@nttbl.jp

¹L. I. Glazman and M. E. Raikh, JETP Lett. **47**, 452 (1988).

²T. K. Ng and P. A. Lee, Phys. Rev. Lett. **61**, 1768 (1988).

³D. Goldhaber-Gordon, H. Shtrikman, D. Mahalu, D. Abusch-Magder, U. Meirav, and M. A. Kastner, Nature (London) **391**, 156 (1998).

⁴S. M. Cronenwett, T. H. Oosterkamp, and L. P. Kouwenhoven, Science **281**, 540 (1998).

⁵J. Schmid, J. Weis, K. Eberl, and K. von Klitzing, Physica B **256B-258B**, 182 (1998).

⁶W. G. van der Wiel, S. De Franceschi, T. Fujisawa, J. M. Elzerman, S. Tarucha, and L. P. Kouwenhoven, Science **289**, 2105 (2000).

⁷A. C. Johnson, C. M. Marcus, M. P. Hanson, and A. C. Gossard, Phys. Rev. Lett. **93**, 106803 (2004).

⁸K. Kobayashi, H. Aikawa, A. Sano, S. Katsumoto, and Y. Iye, Phys. Rev. B **70**, 035319 (2004).

⁹K. Kang, S. Y. Cho, J. J. Kim, and S. C. Shin, Phys. Rev. B **63**, 113304 (2001).

¹⁰M. E. Torio, K. Hallberg, A. H. Ceccatto, and C. R. Proetto, Phys. Rev. B **65**, 085302 (2002).

¹¹I. Maruyama, N. Shibata, and K. Ueda, J. Phys. Soc. Jpn. **73**, 3239 (2004).

¹²M. Sato, H. Aikawa, K. Kobayashi, S. Katsumoto, and Y. Iye, Phys. Rev. Lett. **95**, 066801 (2005).

¹³H. Tamura, K. Shiraishi, and H. Takayanagi, Jpn. J. Appl. Phys., Part 2 **43**, L691 (2004).

¹⁴H. Tamura and L. Glazman, Phys. Rev. B **72**, 121308(R) (2005).

¹⁵N. J. Craig, J. M. Taylor, E. A. Lester, C. M. Marcus, M. P. Hanson, and A. C. Gossard, Science **304**, 565 (2004).

¹⁶P. Simon, Phys. Rev. B **71**, 155319 (2005).

¹⁷P. Simon, R. Lopez, and Y. Oreg, Phys. Rev. Lett. **94**, 086602 (2005).

¹⁸M. G. Vavilov and L. I. Glazman, Phys. Rev. Lett. **94**, 086805 (2005).

¹⁹A. Richter, M. Yamaguchi, T. Akazaki, H. Tamura, and H. Takayanagi, Jpn. J. Appl. Phys., Part 1 **43**, 7144 (2004).

²⁰Transmission line shape is given as $G(\epsilon) \propto (\epsilon+q)^2/(\epsilon^2+1)$, where ϵ is the detuning from the resonance position normalized with the level broadening, and q is Fano's asymmetric parameter.

²¹K. Kobayashi, H. Aikawa, S. Katsumoto, and Y. Iye, Phys. Rev. Lett. **88**, 256806 (2002).

²²D. L. Cox and A. Zawadowski, Adv. Phys. **47**, 599 (1998).

²³G. B. Martins, C. A. Busser, K. A. Al-Hassanieh, E. V. Anda, A. Moreo, and E. Dagotto, Phys. Rev. Lett. **96**, 066802 (2006).